



Efficient 2nd Order Taylor Series Expansion Model to Integrate Stiff and Non-Linear ODEs: Comparison to Rk45-Type Integrators

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Abstract

A Taylor Series Expansion (TaSE) and three Runge-Kutta RK45 embedded pairs numerical integrators of ordinary differential equations (ODEs) are developed. The TaSE model is designed to solve higher order (>1) ODEs. Validation of TaSE was carried out by help of the RK45 models as well as exact solution for well-chosen stiff test cases. The model implementation and its local and global order of accuracy are investigated.

Keywords and Phrases: Taylor Series Expansion model; RK45 integrators; ODEs; Accuracy; Stiffness; Validation

Introduction

Historically, it first started with Ludwig Prandtl in 1904 [1] when he gave a seminar at the International Congress of Mathematics in Heidelberg, under the title "Über Flüssigkeitsbewegungen bei sehr kleiner Reibung". Prandtl took the first steps in a branch of science which turns out to be double-edged; on the one hand for the fact of being branched into several physical areas and on the other hand for the challenges encountered mathematically (existence and uniqueness of the solution) then numerically (solvency, accuracy etc.). Prandtl first initiated the theory of boundary layers problems (BLP) and laid its foundations providing later a great research interest for the several scientific and engineering applications such as the calculation of the drag coefficient of a lifting airfoil or hydrofoil airplane wing, turbine blade [2]. A little bit after, Blasius [3] gave in 1908 his famous non-linear boundary value problem (BVP) which, among others, has been used until today as good test case to the new numerical methods and models of integration of ordinary differential equations (ODEs) used in the following for the resolution of partial differential equations (PDEs).

The investigation of differential equations may be done using the symbolic analytic methods (exact), the geometric methods as well as the numerical methods. Besides, Blasius problem is a BVP, posing mathematical and numerical difficulties to its resolution that should only be numerically. Weyl [4] was the first to prove the existence and uniqueness of the solution of the Blasius problem where an accurate determination of the so-called shear stress $f''(0)$ annoyed researchers. In fact, the Blasius BVP should be transformed to initial value problem (IVP) to complete the explicit and missing equation to the starting closure of the problem. Blasius gave bounds of the value between 0.3315 and 0.33175 and later the value was determined with sufficient precision as $f''(0) = 0.332057336$. using, among other algorithms, the shooting method and its well recommended variants for BVPs [5-12] as well as the finite differences and collocation methods. Parallel to the continuous appearance of new schemes and families of numerical methods for solving ODEs [13-16], ubiquitous stiff problems are emerging namely in electronics, solid mechanics, plasma physics, chemical

kinetics, biology etc. The application of known methods exhibits some instabilities and accuracy loss. To overcome problems of stability, accuracy and ODEs stiffness due to the presence of different time scales in the handled problem it is inevitably to develop new high order flexible and accurate integrators. Many attempts are made to extend the notion of differential equations and the range of available methods. The classification of solution methods to ODEs and systems of ODEs may be performed differently based on some criteria and properties, namely explicit embedded and implicit methods (especially for the Runge-Kutta RK method variants), one-step, two-steps or multisteps methods [17-19].

Over the years, many numerical methods as listed above but of different forms, namely Adams-Bashforth, Adams-Moulton or Multistep, are used on the same principle. Because of its high performance at relatively little computation cost, the RK4 finds great use in wide range of applications even originally has been founded as fixed-step integrator. The RK4 integrator idea holds wide range of interest and gained continuous development [20-23] and application which lead to several improvements such as providing adaptive step size. The latter is a good property that should exert a good integrator over its own time-marching; so that the step size h changes continuously to reduce the computational cost while keeping the total error within given tolerance ϵ . Each 'new' step size is called the optimal step size subject the calculation methods. The integrators using this idea are often called embedded pairs method, such as the RK45 variants: RK-Cash Karp method (RK45-CK), RK-Fehlberg method (RK45-F), RK-Dormand-Prince method (RK45-DP) etc.

However, all numerical methods have their own potential, upsides and downsides when solving various types of ODEs and development of new solution techniques is important. In the present work a Taylor Series Expansion (TaSE designed) model is developed to solve higher order (>1) ODEs and system of ODEs and was compared to three home coded RK45 models as well, namely the RK45-CK, RK45-F and RK45-DP. The models are compared on second and three order ODEs test cases presenting different stiffness scales. The TaSE model is designed to solve higher orders ODEs and system of ODEs and is expected to be at least second order solver since it uses the explicit trapezoidal rule.

Taylor Series Expansion Model: Development and Validation

2nd order B. van der Pol problem

The method is explained by an example. Let's start with the non-linear second order ODE of Balthasar Van der Pol (1926) circuit. A simple model of the ODE is given as follows:

$$f'' + \zeta(f^2 - 1)f' + f = 0 \quad (1)$$

The initial conditions will be chosen as $f(0) = 2$ and $f'(0) = 0$ for the variable η ranging in $[0, 2]$ and $\zeta > 0$ is a parameter.

Note that this ODE may be transformed to a system of ODEs as:

$$\begin{cases} f' = g \\ g' = \zeta(1 - f^2)g - f \end{cases} \quad (2)$$

Even of more higher order, Eq. (1) may be reduced to a similar system of first order ODEs. However, it has been shown that the solution scheme will lose of efficiency.

The Taylor series expansion (TaSE) to second order of f in the vicinity of value η using a small enough $h \ll \eta$, gives:

$$f(\eta+h) = f(\eta) + hf'(\eta) + \frac{h^2}{2!} f''(\eta) + O(h^3) \quad (3)$$

The value of $f''(\eta)$ may be obtained from Eq. (1) as:

$$f(\eta+h) = f(\eta) + hf'(\eta) + \frac{h^2}{2!} (\zeta(1 - f^2(\eta))f'(\eta) - f(\eta)) + O(h^3) \quad (4)$$

To calculate $f(\eta+2h)$ (the second iteration), the value of $f'(\eta+h)$ should be computed before. The value may be obtained differently by using a first order TaSE as:

$$f'(\eta+h) = f'(\eta) + hf''(\eta) + O(h^2) \quad (5)$$

or by integration of Eq. (1) as:

$$\int_{\eta}^{\eta+h} f''(t) dt = \int_{\eta}^{\eta+h} (\zeta(1 - f^2) f' - f) dt \quad (6)$$

$$f'(\eta+h) - f'(\eta) = \zeta(1 - f^2(t_{\eta})) \int_{\eta}^{\eta+h} f'(t) dt - \int_{\eta}^{\eta+h} f(t) dt \quad (7)$$

Where $t_{\eta} \in [\eta+h, \eta]$ and the quantity $(1 - f^2(t_{\eta}))$ comes out of the integral using the theorem of the mean value (implicitly),

$$\int_{\eta}^{\eta+h} \phi(t) f'(t) dt = \phi(t_{\eta}) \int_{\eta}^{\eta+h} f'(t) dt \quad (8)$$

and can be estimated as $\phi(t_{\eta}) \approx (\phi(\eta+h) + \phi(\eta)) / 2$ and the integration of f at the last right hand side can be estimated to second order using the trapezoidal rule as $h(f(\eta+h) + f(\eta)) / 2 - h^3 / 12 f''(\xi_n)$. However, the integration of f' is exact, which reduces the truncation error as:

$$f'(\eta+h) - f'(\eta) = \zeta(1 - f^2(t_{\eta})) (f'(\eta+h) - f'(\eta)) - \frac{h}{2} (f(\eta+h) + f(\eta)) + O(h^3) \quad (9)$$

Arranging terms leads to the following formula for f' :

$$\underbrace{f'(\eta+h) - f'(\eta)}_{\text{Exact}} = \underbrace{\zeta(1 - f^2(\eta+h/2))}_{+O(h^3)} \underbrace{(f(\eta+h) - f(\eta))}_{\text{Exact}} - \underbrace{hf(\eta+h/2)}_{O(h^3)} + O(h^3) \quad (10)$$

One may also calculate $f'(\eta+h)$ using first order scheme using the increase rate of f or the slope f' . Note, also, that in the present case the first right term of Eq. (7) may be integrated directly since the simple form of $(1 - f^2) f'$. The needed quantities for the calculation of $f(\eta+h)$ at Eq. (3) are determined. It is expected to have here a second order accurate solution.

The comparison of the TaSE results against RK45-F and RK45CK ones for the ζ parameter values 0.5 and 5 are plotted in Figures 1 and 2.

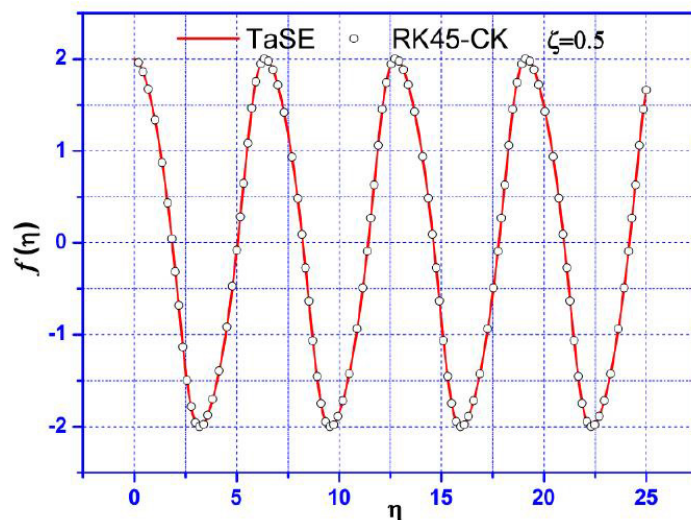


Figure 1: Comparison between the results of the present models TaSE and RK45-CK for $\zeta = 0.5$. Initial step size.

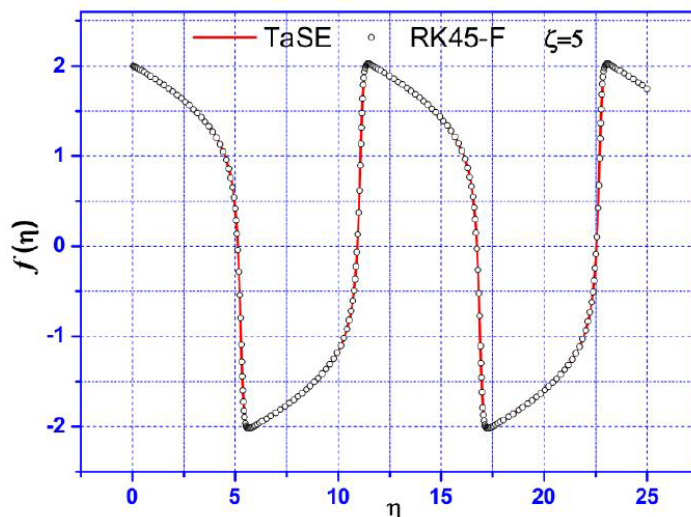


Figure 2: Comparison between the results of the present models TaSE and RK5-F for $\zeta = 5$. Initial stepsize $h=0.2$.

The model's results are in excellent agreement. The comparison case $\zeta = 5$ reveals the problem of flame propagation whose simple model is described by the

$$\frac{dy}{dt} = y^2 - y^3, y(0) = \delta, 0 \leq t \leq \frac{2}{\delta} \quad (11)$$

The problem is very stiff for very low values of the parameters δ called the initial radius. An investigated method should exhibit sufficient aptitude to simulate the very acute jump of the flame radius. In Figure 2, the two models seem to be efficient to handle such problems.

The numerical results of comparison of the TaSE and RK45-F models are given in Table 1. By checking the order of the TaSE

model (local error not global) gives that is of second order accuracy. Verification can be performed on the results for the ability to capture a peak, say for example for $\eta = 7$ and for different step sizes $h = 0.0005, h = 0.005$ and $h = 0.05$. Assuming the truncation error $\Delta f \sim O(h^a) = bh^a$. The plot of $\ln(\Delta f)$ vs. $\ln(h)$ is given in

Figure 3. The fit gives $\ln(\Delta f) = 1.975 \ln(h) + 0.95$, so the local deviation under the prescribed conditions ($\zeta = 5, \eta = 7$) is $\Delta f \approx 2.4647h^{1.966}$. The fit was obtained within $R^2 = 0.99965$ confidence.

Note that the results of Table 1 are obtained for $\zeta = 5$, TaSE model: $h = 0.5 \cdot 10^{-3}$, RK45-F and RK45-CK: fixed step $h = 10^{-3}$ and $\epsilon = 10^{-5}$ and that the RK45-F integrator is found to be slightly more fast than RK45-CK one regarding the extra number of iterations taken sometimes for adaptive stepsize.

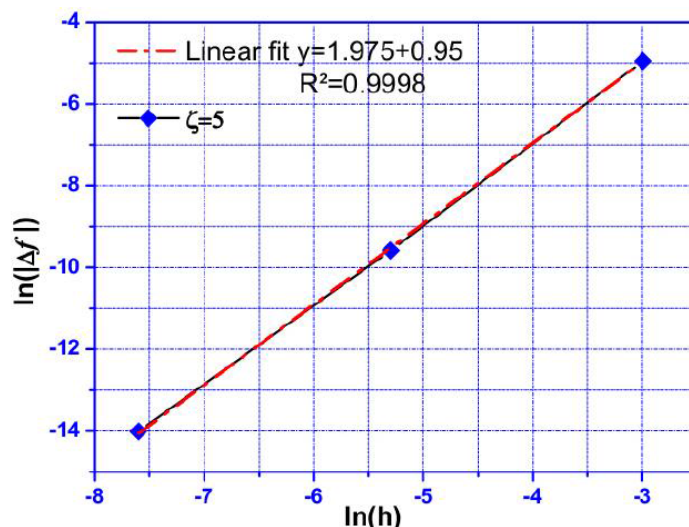


Figure 3: Estimation of the TaSE model order of accuracy.

Table 1: Comparison of the present results for three models: TaSE, RK45-F and RK45-CK with relative deviation. The deviation is calculated between TaSE and RK45-F results.

Present Models Result for f				
η	TaSE	RK 45-F	RK45-CK	Error
0.5	1.941176184	1.941176665	1.94117689	-1.27E-07
1	1.869438854	1.86943948	1.86943972	2.51E-03
1.5	1.792576327	1.79257715	1.79257739	-1.05E-07
2	1.709233873	1.70923531	1.70923555	-5.89E-07
2.5	1.6173270 12	1.61732924	1.61732948	-1.55542E-07
3	1.5133890 13	1.51339197	1.5133922 1	-4.97955E-07
3.5	1.390924524	1.39092922	1.39092946	5.47042E-07
4	1.235281558	1.23528850	1.23528862	-9.24522E-07
4.5	1.000752540	1.00076473	1.00076449	-L7604E-06
5	0.404182268	0.40422821	0.404221	-7.34448E-06
5.5	-1.971318608	-1.97129858	-1.97128713	-2.61699E-06
6	-1.982958921	-1.98296154	-1.98295105	9.62396E-07
6.5	-1.91390362	-1.91390646	-1.91389549	5.02542E-07
7	-1.840305439	-1.84030461	-1.8402971	8.26161E-07
7.5	-1.761122436	-1.76112616	-1.76111364	8 0526E-07

One can observe in the vector field graph (x, x') of Figure 2, the so-called Poincaré limit cycle which depends strongly on the ζ value. For instance, low ζ values ($\zeta \ll 1$) lead to circular limit cycle since Eq. (1) becomes equivalent to harmonic equation. The graph solution tends to a circle as shows Figure 4.

The excellent agreement obtained in the first test case shows that the results of the TaSE model exhibit high efficiency in joining integrators of high performance such embedded pairs RK models.

3rd order Blasius and Sakiadis problems

The Blasius problem is a third order nonlinear BVP with single ODE. The well-known ODE is given as:

$$f''' + \frac{1}{2} f f'' = 0 \quad (12)$$

With the conditions $f(0) = f'(0) = 0$ and $f''(\infty) = 1$ for the Blasius problem and $f(0) = f''(\infty) = 0$ and $f'(0) = 1$ for the Sakiadis problem, which are two BVPs. The two should be transformed to IVP to be solved.

In the following the Blasius problem will be investigated. As was shown in the first example, the equations used for simulations are:

$$f(\eta+h) = f(\eta) + hf'(\eta) + \frac{h^2}{2!} f''(\eta) + \frac{h^3}{3!} \left(-\frac{1}{2} f(\eta) f''(\eta) \right) + O(h^4) \quad (13a)$$

$$f'(\eta+h) = f'(\eta) + hf''(\eta) + \frac{h^2}{2!} \left(-\frac{1}{2} f(\eta) f''(\eta) \right)' + O(h^3) \quad (13b)$$

$$f''(\eta+h) = f''(\eta) - \frac{1}{4}(f(\eta+h) + f(\eta))f'(\eta+h) - f'(\eta) + O(h^3) \quad (13c)$$

The obtained results for TaSE and RK45-F models gathered with literature results are given in Table 2. The results of the proposed TaSE model are in excellent agreement with those of the

RK45-F integrator either for f and f' even for quite coarse step-size $h = 0.01$; which reduces considerably the number of iterations. The computed quadratic (RMS) deviation is close to $1.35554E-05$ for f and $6.17489E-06$ for f' which supports well on the accuracy of the proposed model.

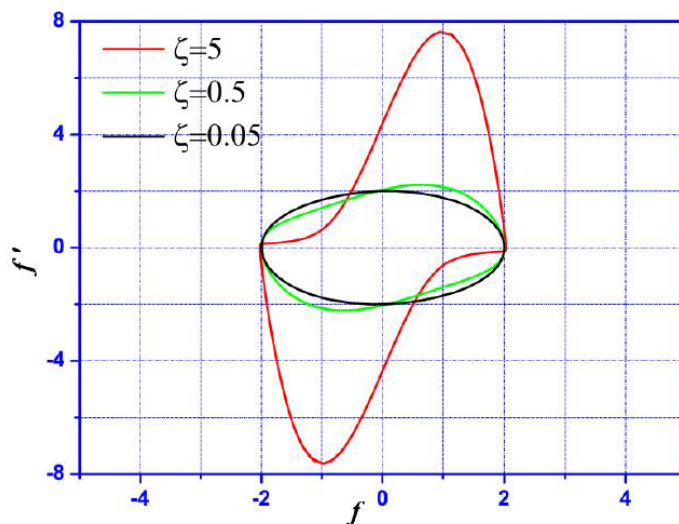


Figure 4: Poincare limit cycle for $\zeta = 0.5$ and $\zeta = 5$. The black cycle is x-elongated due to the $f - f'$ scale difference.

Table 2: Comparison of the present results for the models TaSE and RK45-F with reference ones for f and f' . The present results are obtained for $h=0.01$ for TaSE and $h=0.001$ for RK45-F.

η	Results for f			Results for f'		
	TaSE	RK45-F	Ref. [24]	TaSE	RK45-F	Ref. [25]
1	0.165571681	0.16557156	0.1656	0.329780011	0.32977977	0.32978
2	0.650024365	0.65002447	0.65	0.629765774	0.62976617	0.62976
3	1.396807934	1.39681005	1.3969	0.846043622	0.84604508	0.84604
4	2.305744552	2.30574822	2.3059	0.955515949	0.95551944	0.95551
5	3.283269 106	3.28327489	3.2834	0.991538971	0.99154246	0.99154
6	4.279613436	4.27962494	4.2798	0.99896999	0.99897337	0.99897

System of coupled 2nd order ODEs

We consider the following oscillatory nonlinear system of second order ODEs with initial values:

$$x'' = -4t^2 x - \frac{2y}{\sqrt{x^2 + y^2}}, \quad x(0)=1, \quad x'(0)=0 \quad (14a)$$

$$y'' = -4t^2 y - \frac{2x}{\sqrt{x^2 + y^2}}, \quad y(0)=1, \quad y'(0)=0 \quad (14b)$$

Whose analytic solution represents a periodic motion with variable period and is given by $x(t) = \cos(t^2)$ and $y(t) = \sin(t^2)$.

As was shown here above examples, the equations used for the simulation of x for example are:

$$x(t+h) = x(t) + hx'(t) + \frac{h^2}{2!} x''(t) + O(h^3) \quad (15a)$$

$$x'(t+h) = x'(t) + \left(-4t^2 y + \frac{2x}{\sqrt{x^2 + y^2}} \right)_{t=t+h/2} + O(h^3) \quad (15b)$$

$$x''(t+h) = x''(t) + \int_t^{t+h} \eta^2 d\eta - \frac{1}{2} f(t+h/2) + O(h^3) \quad (15c)$$

Where

$$f(t) = \frac{2y(t)}{\sqrt{x(t)^2 + y(t)^2}} \quad (16)$$

The present example shows how to estimate $x'(t+h)$ to second order in a second order ODE. The problem is solved for $t \in [0, 7]$ with step size $h = 0.001$.

The results of TaSE and RK45-DP models are gathered in Table 3 with their deviations from the exact solution. Extra digits are given here since the comparison is performed against exact

solution. The plot of the two computational results is depicted in Figure 5. Results of Table 3 and Figure 5 reveal the powerfulness of the TaSE model. Besides, in this test case we have reviewed the global order of accuracy of the TaSE model. The RMS deviation over

the seven computed times was calculated for three stepsizes: $h = 0.001$, $h = 0.005$ and $h = 0.01$. It has been found that $\ln(\Delta f) = 1.998 \ln(h) + 4.445$ with a coefficient of determination $R^2 = 1$. The expected order of accuracy is always preserved.

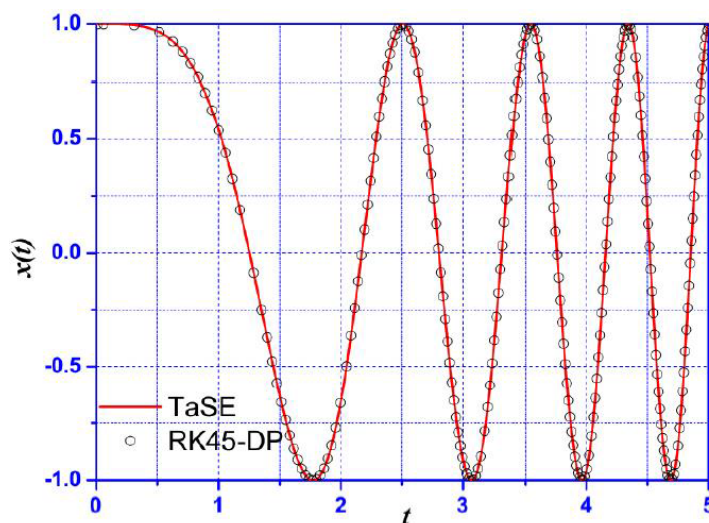


Figure 5: Plots comparing TaSTE and RK45-DP numerical results for a system of coupled 2nd order ODEs.

Table 3: Comparison of the present results for the models TaSE and RK45-DP with exact solution for a step size $h=0.001$ for both.

t	TaSE	RK45-DP	Err 1	Err 2
1	0.54030228	0.540302306	4.26E-07	-2.59E-08
2	-0.65364355	-0.653643621	-9.52E-07	7.10E-08
3	-0.911130352	-0.911130262	-8.44E-06	-9.05E-08
4	-0.957659358	-0.95765948	-4.65E-06	1.22E-07
5	0.991202835	0.991202812	2.42E-05	2.30E-08
6	-0.127962832	-0.12796369	1.04E-4	8.58E-07
7	0.300593577	0.300592544	2.01E-04	1.03E-06

Conclusion

A Taylor Series Expansion numerical integrator of ODEs is developed. The model is exerted to solve higher order ODEs. The following concluding remarks may be drawn:

- The model has been validated on three fourth order accurate RK models and exact solution and gave excellent agreement for three test cases of second and three order ODEs of different stiffness scales.
- The local and global order of accuracy of the model was checked and estimated to be second order accurate.
- The model is always efficient even for first approximation of the first derivative for second order ODEs.
- The Trapezoidal explicit rule and the theorem of the mean value are used when needed to handle to second order very complex and hard-to-derive terms.

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None.

Conflict of Interest

No conflict of interest.

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